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LOCAL STATES AND MAGNETIC INTERACTIONS STUDY OF TRANSITION METAL IONS IN VITREOUS OXIDE MATRICES

PhD Thesis Summary

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INTRODUCTION

Glasses are noncrystaline solid materials that have a disordered structure, but have a local order. The study of materials with vitreous structures gained a special attention due to their manifold applications in domains of science and technology. The oxide glasses unfold a large variety of properties due to the variety of their compounds; these glasses have wide technological applications.

This thesis work is focused on characterizing the structure, the optic and magnetical properties of arsenolite and borate glasses. The study presented in this thesis has as objective the obtaining of new experimental data and the clarification of the physical phenomena which determine the structural behavior of As₂O₃ and B₂O₃ based glasses doped with manganese or iron ions. In order to obtain this information of $xMO \cdot (100-x)[As_2O_3 \cdot TeO_2]$, $xMO \cdot (100-x)[As_2O_3 \cdot TeO_2]$ and $xMO \cdot (100-x)[3B_2O_3 \cdot Li_2O]$, where MO => MnO or Fe₂O₃ glass systems, with 0 < x < 50% mol MnO or Fe₂O₃, they were prepared and investigated by x - ray Diffraction, FT - IR spectroscopy, Raman spectroscopy, electron paramagnetic resonance (EPR) and magnetic susceptibility measurements.

The thesis is structured in four chapters. The first chapter presents, on literature basis, the main results regarding the structure and the properties of As_2O_3 and B_2O_3 based glasses.

Chapter 2 presents theoretical and experimental general aspects regarding the iron and manganese behavior using the IR absorption, Raman spectroscopies, electron paramagnetic resonance and magnetic susceptibility. Chapter 3 describes the experimental techniques used in the studies of glasses. The processing and preparation of the samples are presented. Chapter 4 presents and discusses the experimental results obtained for the investigated glasses. In the end, the general conclusions which emphasize the most important results of this thesis are presented.

CHAPTER 1. THE STUDY OF THE STRUCTURE OF As₂O₃ AND B₂O₃ BASED GLASSES

 As_2O_3 and B_2O_3 are the most important glass forming oxides and have been incorporated into various kinds of glass system in order to obtain the desired physical and chemical properties for both scientific and industrial applications. As_2O_3 is a very strong glass formation having a configuration based on pyramidal AsO_3 units, determined by a special dynamics on the As and O bonds, unique properties of the glass systems [1,-3]. The properties of the glasses based the B_2O_3 are due to the fact that the boron atom can assume triangular and tetrahedral coordination and also to the different ways through which the borate buildings units can be linked together [4-9].

The behavior of various borate species in the glass structure is determined by the nature of modifier oxides (TeO₂, PbO, etc.). In most binary borate glasses the glass properties change almost linearly with composition [10-17]. In binary borate glasses the properties of these show a major deviation from linearity which depends on modifier oxide. The structural reasons for this is generally due to the change in the coordination of boron atom from 3 to 4 and back again. Many technological applications of these glasses depend of this effect since is produces many improvements [18-23].

The introduction of MnO [24-32] or Fe_2O_3 [33-45] into diamagnetic glasses transform them into paramagnetic glasses. The introduction of iron and manganese in small percentages modifies the physical properties of the glasses and the material is more accessible for a complete structural characterization.

In order to extend the available information concerning the interesting class of As_2O_3 and B_2O_3 based glasses we investigated glasses by means of different physical methods. We obtained valuable information concerning the structure and structural changes of the glasses that occur with the addition of transitional oxide metals.

CHAPTER 2. METHODS USED IN THE STUDY OF STRUCTURE AND PROPERTIES OF OXIDE GLASSES

In this chapter some of the methods used in the investigation of the structure and properties of vitreous systems are presented. The study of the structure and properties of oxide glasses with transitional metals ions: IR and Raman spectroscopies, electron paramagnetic resonance and magnetic susceptibility measurements were chosen as investigation methods.

The electromagnetic radiation (EM) with frequencies between 4000 and 400 cm⁻¹, called infrared radiation (IR), is used to obtain information connected to the structure of the compound and as an analytic method to test the purity of the compound. Its application in organic chemistry is known as IR spectroscopy. The radiation is absorbed by organic molecules and converted in molecular vibration energy. When the radiant energy coincides with a specific molecular vibration the absorption appears. The wave length at which the radiation is absorbed by the organic molecule gives information regarding the functional groups presented in the molecule [1-4].

The Raman spectroscopy brings useful contributions to the molecular vibrations study. The interraction mechanism of electromagnetic radiation with the molecular vibration which forms the Raman spectroscopy theoretical basis differs from the one of the process which forms the IR spectroscopy theoretical basis [4]. Another method used in the investigation of the glass system is the electron paramagnetic resonance (RPE). This method is one of the most powerful techniques for local order investigation. RPE is a method widely used to describe the fundamental states and to characterize the vicinities effect on the energetic levels of the paramagnetic centers. The method consists in the study of the electronic split levels of the atoms in the presence of an external magnetic field [6]. The Mn²⁺ EPR spectra are characterized by resonance absorptions at $g_{ef} \approx 4.3$, $g_{ef} \approx 3.3$ and $g_{ef} \approx 2.0$. The resonance line centered at $g_{ef} \approx 4.3$ is corresponding to the isolated Mn²⁺ ions. The line from $g_{ef} \approx 2.0$ is attributed to Mn²⁺ ions involved in magnetic interactions.

The Fe³⁺ EPR spectra are characterized by resonance absorptions at $g \approx 9,7$, $g \approx 6$, $g \approx 4,3$ and $g \approx 2,0$. The resonance line from $g \approx 9,7$ is related to Fe³⁺ ions disposed in rhombic environment. The resonance line at $g \approx 6$ was assigned to an axial distortion of the paramagnetic ions neighborhood. The resonance line at $g \approx 4,3$ is corresponding to the isolated Fe³⁺ ions situated in octahedral, rhombic or tetragonal symmetric distorted neighborhoods. The line from $g_{ef} \approx 2,0$ is attributed to Fe³⁺ ions involved in magnetic interactions.

Along with the EPR spectroscopy, the magnetic susceptibility measurements provide useful information about the valence state and the interaction involving the transitional metal ions in vitreous materials [7-10].

Different vitreous systems with manganese or iron ions were intensively studied by magnetic susceptibility measurements to determine the magnetic behavior and the valence states of these ions in these systems. It was observed that the presence of the manganese or iron ions in different valence states depends on the chemical composition of the vitreous matrix, on the valence of the vitreous network former and modifications and on the preparation of the environment. The magnetic properties of the oxide glasses with manganese or iron ions were assigned to the antiferromagnetic couple between Mn²⁺-Mn²⁺, Mn²⁺-Mn³⁺ and Mn³⁺-Mn³⁺ [19-26] or Fe -Fe³⁺, Fe²⁺-Fe²⁺ and Fe³⁺-Fe²⁺ ions [11-18].

CHAPTER 3. EXPERIMENTAL TECHNIQUES

We have prepared glasses of the $xMO \cdot (100-x)[As_2O_3 \cdot TeO_2]$, $xMO \cdot (100-x)[As_2O_3 \cdot PbO]$ şi $xMO \cdot (100-x)[3B_2O_3 \cdot Li_2O]$ glass structure, where MO => MnO or Fe_2O_3 with 0 < x < 50 % mol, used components of reagent grade purity: $MnCO_3$, Fe_2O_3 , As_2O_3 , H_3BO_3 , PbO, Li_2CO_3 şi TeO_2 in suitable proportions to obtain the desired composition. The mixtures were melted in sintered corundum crucibles, introduced in an electric furnace directly at 1250 C and kept for 5 minutes at this temperature. They were quickly cooled at room temperature by pouring onto stainless steel plates.

The structure of the samples was analyzed by means of x-ray Diffraction, using powders, with a Bruker D8 Advanced diffract meter.

The FT - IR spectra have been recorded using a Bruker Equinox 55 with a spectral range between 4000 cm⁻¹ and 370 cm⁻¹. A MIR, GLOBAR generator cooled with air was used. The detection was carried out with a DLATGS detector with a KBr window. The spectral resolution was about 0.5 cm^{"1}. The samples were prepared using KBr pellet technique.

The Raman spectra were recorded, on bulk samples, at room temperature using an integrated FRA 106/S Raman module attached to Bruker Equinox 55 with a spectral range from 3600 cm⁻¹ to 70 cm⁻¹. An Nd.YAG laser with an output power of 500 mW and a 1064 nm radiation was used. The detection was carried out with an ultra sensitive D418-T detector cooled with liquid nitrogen. The spectral resolution was about 1 cm⁻¹.

The EPR spectra were obtained at room temperature with an Adani Portable EPR Spectrometer PS8400 in X-frequency band (9.4 GHz). For these measurements, equal quantities of powders from the investigated samples, closed in glass tubes, were used.

Magnetic susceptibility measurements were performed on a Faraday type balance in the 80 - 300 K temperature range. The sensitivity of the equipment was 10^{-7} emu/g.

CHAPTER 4. RESULTS AND DISCUSSION REGARDING THE STUDY OF xMO·(100-x)[As₂O₃·TeO₂], xMO·(100- x)[As₂O₃·PbO] şi xMO·(100-x)[3B₂O₃·Li₂O] GLASS STRUCTURE, WHERE MO => MnO or Fe₂O₃

In this chapter the results obtained by following the study of the structure and properties of oxide glasses with transitional metals ions by FT- IR and Raman spectroscopies, electron paramagnetic resonance and magnetic susceptibility measurements are presented. The structural changes and magnetic properties with the manganese and iron content were followed.

4.1. Comparative study by FT-IR and Raman spectroscopies of xMO·(100x)[As₂O₃·TeO₂], xMO·(100- x)[As₂O₃·PbO] şi xMO·(100-x)[3B₂O₃·Li₂O] glasses, where MO => MnO or Fe₂O₃

4.1.1. Comparative study by FT-IR spectroscopy of As₂O₃·TeO₂ and As₂O₃·PbO vitreous matrices

In order to obtain new information regarding the role of the modifier in the formation of vitreous matrices, comparative studies by FT-IR spectroscopy were made. The FT-IR spectroscopy is represented in figure 4.1(a,b). The wave length and the structural assignments of FT – IR spectra of As₂O₃·PbO and As₂O₃·TeO₂ vitreous matrices are represented in table 4.1.

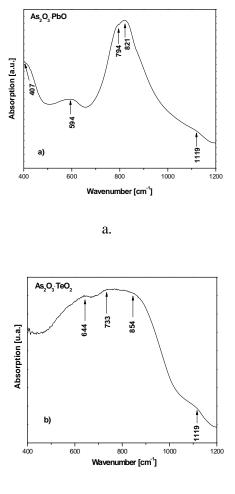




Fig.4.1. FT – IR spectra of vitreous matrices, a. As₂O₃·PbO și b. As₂O₃·TeO₂.

The FT – IR spectrum of As_2O_3 ·PbO glass matrix presents five absorption bands at: ~ 1119 cm⁻¹, ~ 821 cm⁻¹, ~ 794 cm⁻¹, ~ 594 cm⁻¹ and ~ 407 cm⁻¹(Fig.4.1.a).

The band from ~ 1119 cm⁻¹ is due to totally symmetric stretching vibrations of AsO₃ units [5 - 7]. The bands from ~ 821 cm⁻¹ and ~ 794 cm⁻¹ are due to doubly degenerate stretching vibrations of AsO₃ structural units [5 - 7]. The large band centered at ~ 594 cm⁻¹ can be due to totally symmetric bending vibrations of AsO₃ structural units [5-7]. The band from ~ 407 cm⁻¹ can be due to vibrations of Pb – O bonds from PbO₄ structural groups [7].

Following As_2O_3 TeO₂ glass matrix spectrum presented in Fig. 4.1.b, four bands can be observed located at: ~ 644 cm⁻¹, ~ 733 cm⁻¹, ~ 854 cm⁻¹ and ~ 1119 cm⁻¹. The band from ~ 644 cm⁻¹ can be assigned to totally symmetric bending vibrations of AsO₃ units [5-7] and also stretching vibration mode of TeO₄ tbp with bridging oxygen[11]. The band from ~ 733 cm⁻¹ can be assigned to the symmetric stretching vibration mode of TeO bonds in TeO₃ units [15]. The band from ~ 854 cm⁻¹ can be assigned to doubly degenerate stretching vibrations of AsO₃ units [5-7] and also to the stretching vibration mode of TeO₃ tp with NBO [11]. The band from ~ 1119 cm⁻¹ can be assigned to totally symmetric stretching vibrations of AsO₃ units [5-7]. These bands, presented in As₂O₃·TeO₂ glass matrix spectrum, confirm the presence of AsO₃, TeO₄ tbp and TeO₃ tp units in the structure of the glass matrix but, due to the broadness of the bands, it can not be told for sure which one of these units is predominant.

$\widetilde{\upsilon}$ [cm ⁻¹]		Attribution	
As ₂ O ₃ ·PbO	As ₂ O ₃ ·TeO ₂	Autouton	
~1119 cm ⁻¹	~ 1119 cm ⁻¹	Totally symmetric stretching vibrations of AsO ₃ units	
~ 821cm ⁻¹		Doubly degenerate stretching vibrations of AsO ₃ structural units	
	~ 854 cm ⁻¹	Doubly degenerate stretching vibrations of AsO ₃ units Stretching vibration mode of TeO ₃ tp with NBO	
~ 794 cm ⁻¹		Doubly degenerate stretching vibrations of AsO ₃ structural units	
	~ 733 cm ⁻¹	Symmetric stretching vibration mode of Te-O bonds in TeO ₃ units	
~ 594 cm ⁻¹	~ 644 cm ⁻¹	Totally symmetric bending vibrations of AsO ₃ units	
$\sim 407 \text{ cm}^{-1}$		Vibrations of Pb – O bonds from PbO ₄ structural groups	

Table 4.1.Wavenumber and the structural assignments of FT – IR spectra of As₂O₃·PbO and As₂O₃·TeO₂ vitreous matrices.

4.1.2. Comparative study by FT-IR spectroscopy of xMO·(100- x)[As₂O₃·PbO] vitreous systems, where MO => MnO or Fe₂O₃

The infrared absorption spectra obtained for $xMO(1-x)[As_2O_3 PbO]$ glass system, where MO => MnO or Fe_2O_3 , with $0 \le x \le 50$ mol%, are presented in figure 4.2-3 and their structural assignments are summarized in table 4.1.

The spectra were discussed on the basis of the method given by Tarte [3] and Condrate [4] by comparing the experimental data of glasses with those of related crystalline compounds. The characteristic absorption bands for crystalline As_2O_3 [5-8], PbO [7,9] and characteristic bands for Fe₂O₃ [18,19] were used as a reference point in the results discussion.

The FT – IR spectra of the investigated glasses suggest a structure formed of AsO₃ and PbO₄ structural units where As₂O₃ plays the role of the glass former and PbO plays the role of the glass modifier. The evolution of the spectra with the addition and the increasing of the MnO content suggest that the manganese ions break up a part of As – O – As, Pb – O – Pb and probably of As – O – Pb bonds. The addition of MnO in the As₂O₃·PbO glass matrix is leading to a disordering of the glasses structure with the increasing of the MnO content.

The evolution of the absorption bands for FT - IR spectrum of $xFe_2O_3 \cdot (1-x)$ [As₂O₃·PbO] glasses are determined by the addition of Fe₂O₃ content. It can be observed to shift the bands assigned As - O - As vibrations from AsO_3 units due to changing to length As - O - As bonds and interne angles of AsO_3 units, disappears PbO₄ units and appears FeO₆ units The addition of Fe₂O₃ in the As_2O_3 ·PbO glass matrix determined a disordering of the glasses structure with the increasing of the Fe₂O₃ content.

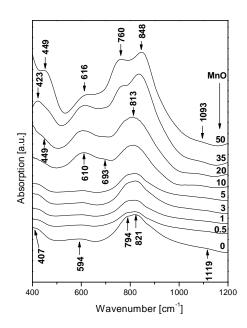


Fig. 4.2. FT – IR spectra of xMnO·(100- x)[As₂O₃·PbO] glasses.

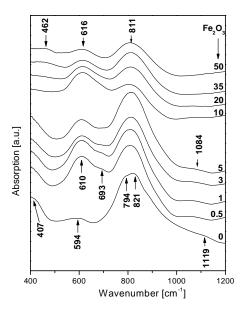


Fig.4.3. FT – IR spectra of xFe₂O₃·(100- x)[As₂O₃·PbO] glasses.

4.1.3. Comparative study by FT-IR spectroscopy of xMO·(100-x)[As₂O₃·TeO₂] vitreous systems, where MO => MnO or Fe₂O

The infrared absorption spectra obtained for $xMO \cdot (1-x)[As_2O_3 \cdot TeO_2]$ glass system, where MO => MnO or Fe_2O_3 , with $0 \le x \le 50$ mol%, are presented in figure 4.4 -5 and their structural assignments are summarized in table 4.1.[16,17]. The characteristic absorption bands for crystalline As_2O_3 [5-8], PbO [7,9], TeO₂ [8,10-15], MnO₂ and Mn₃O₄ [9] and characteristic bands for Fe_2O_3 [18,19] were used as a reference point in the results discussion.

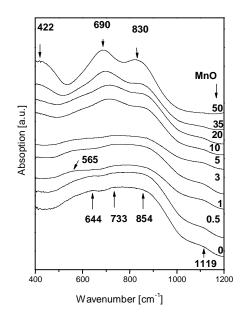


Fig. 4.4. FT – IR spectra of xMnO·(100- x)[As₂O₃·TeO₂] glasses.

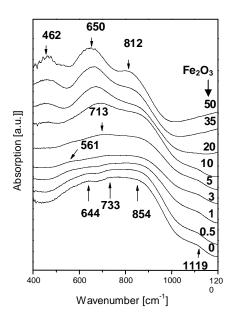


Fig.4.5. FT – IR spectra of xFe₂O₃·(100- x)[As₂O₃·TeO₂] glasses.

This behavior of the studied glasses with the addition and increasing of MnO content suggests that the structure observed in the glass matrix is conserving up to 5 mol% MnO, where the structure takes another form. From a structure dominated by AsO3, TeO4 tbp and TeO₃ tp units, the structure became dominated by TeO₄ tbp and MnO_2 units and in a small measure of AsO_3 and TeO_3 tp units. For higher concentrations, it can not be specified exactly which of the last two structural units have the biggest proportion in these glasses. From FT-IR absorption spectra of the investigated glass system it can be observed that the addition and increasing of manganese ions concluded: the contraction of As-O-As bonds from totally symmetric stretching vibrations of AsO₃ units; for $x \ge 10$ % mol there can be an increase in number of As-O-As bonds from AsO₃ units or/and a length these linkage to doubly degenerate stretching vibrations of AsO₃ units and a stretching vibration mode of TeO₃ tp with NBO. It can be remarked from the FT-IR absorption spectra, the gradual transformation of symmetric stretching vibration mode of Te-O bonds from TeO_3 units in TeO_4 (tbp) units and the presence of MnO₂ units. Once, with the addition of iron ions, the $xFe_2O_3 \cdot (100-x)[As_2O_3 \cdot TeO_2]$ glasses behave like $xMnO(100-x)[As_2O_3 TeO_2]$ glasses, but in this system the appearance of FeO_6 units can be seen.

4.1.4. Comparative study by FT-IR spectroscopy of xMnO·(100- x)[As₂O₃·BO] vitreous systems, were BO => TeO₂ or PbO

The infrared absorption spectra obtained for xMnO·(100- x)[As₂O₃·BO] vitreous systems, where BO => TeO₂ or PbO, with $0 \le x \le 50$ mol% MnO, are presented in figure 4.2 and 4.4. Following the infrared absorption spectra, it can be observed that, for xMnO·(100- x)[As₂O₃· PbO] glass system the most disordered structure to exist for 0 $\le x \le 5$ mol% MnO and to xMnO·(100-x)[As₂O₃· TeO₂] glass system for $1 \le x \le 5$ mol% MnO. For $x \ge 5$ % mol MnO, manganese ions break up a part of As – O – As, Pb – O – Pb and probably of As – O – Pb bonds, to favor the increase in number of PbO₄ units and appearance of MnO₂ units. For $x \ge 10$ % mol MnO, observed comparatively in both system glasses, the addition of manganese ions suggests for the first system a modification in linkage characteristic units and for the second system an improvement for structural units of As₂O₃· TeO₂ vitreous matrice also linkage characteristic units.

4.1.5. Comparative study by FT-IR spectroscopy of $xFe_2O_3(100-x)[As_2O_3BO]$ vitreous systems, where BO => TeO₂ or PbO

The infrared absorption spectra obtained for $xFe_2O_3 \cdot (100-x)[As_2O_3 \cdot BO]$ vitreous systems, where BO => TeO₂ or PbO, with $0 \le x \le 50$ mol% Fe₂O₃, are presented in figure 4.3 and 4.5.

The evolution of the spectra with the addition and the increasing of the Fe₂O₃ content suggest that the iron ions break up a part of As – O – As, Pb – O – Pb and probably of As – O – Pb bonds. It can be observed a shift in the bands assigned As – O – As vibrations from AsO₃ units due to changing to length As – O – As bonds and interne angles of AsO₃ units, the PbO₄ units disappear and FeO₆ units appear. For $x \ge 5$ % mol Fe₂O₃, observed comparatively in both system glasses, the addition of iron ions suggests for the first system a modification in linkage characteristic units of As₂O₃·PbO and for the second system an improvement for structural units of As₂O₃· TeO₂ vitreous matrice also linkage characteristic units, appearance of FeO₆ units in both system.

4.1.6. Comparative study by Raman spectroscopy of As₂O₃·TeO₂ and As₂O₃·PbO vitreous matrices

In order to obtain new information regarding the role of the modifier in the formation of vitreous matrices, there were made comparative studies by Raman spectroscopy. The Raman spectroscopy are presented in figure 4.6(a,b). The wave length and the structural assignments of Raman spectra of As_2O_3 ·PbO şi As_2O_3 ·TeO₂ vitreous matrices are presented in table 4.2.

Table 4.2.Wavenumber and the structural assignements of Raman spectraof As2O3·PbO and As2O3·TeO2vitreous matrices.

$\widetilde{\upsilon} [\mathrm{cm}^{-1}]$		A theilention	
As ₂ O ₃ ·PbO	As ₂ O ₃ ·TeO ₂	- Attribution	
~826 cm ⁻¹	~ 816 cm ⁻¹	Stretching vibrations of As-O-As	
	~ 752 cm ⁻¹	Vibrations of the continuous TeO ₄ (tbp) network	
~ 627 cm^{-1}		Vibrations As – O – As bonds from As ₂ O units	
	~ 580 cm ⁻¹	Stretching vibrations of As-O-As	
$\sim 525 \text{ cm}^{-1}$		Vibrations O-As-O bonds from AsO ₃ units	
	~ 438cm ⁻¹	Symmetric bending vibrations of $Te - O - Te$ vertex linkages between different TeO_4 (tbp), TeO_{3+1} polyhedra Stretching vibrations of As-O-As of AsO ₃ units	
~ 345 cm ⁻¹		Bending vibrations of $As - O - As$ from As_2O units	

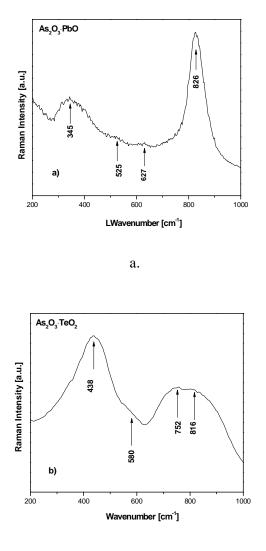




Fig.4.6. Raman spectra of vitreous matrices, a. As₂O₃·PbO și b. As₂O₃·TeO₂.

The Raman spectrum of As_2O_3 ·PbO glass matrix presents four bands at: ~ 826 cm⁻¹, ~ 627 cm⁻¹, ~ 525 cm⁻¹ and ~ 345 cm⁻¹. The band from ~ 826 cm⁻¹ can be assigned to As – O – As stretching vibrations [20]. The band centered at ~ 627 cm⁻¹ can be assigned to vibrations As – O – As bonds from As₂O units, the band from ~ 525 cm⁻¹ are assigned vibrations O-As-O bonds from AsO₃ units [21]. The band of ~ 345 cm⁻¹ are attributed to As – O – As bending vibrations from As₂O units[20,21].

The Raman spectrum of As₂O₃· TeO₂ glass matrix presents four bands at: ~ 816 cm⁻¹, ~ 752 cm⁻¹, ~ 580 cm⁻¹ and ~ 438 cm⁻¹. The band from ~ 816 cm⁻¹ can be assigned to As – O – As stretching vibrations [20]. The band centered at ~ 752 cm⁻¹ can be assigned to vibrations of the continuous TeO₄ (tbp) network [22]. The band of ~ 580 cm⁻¹ are attributed to As – O – As stretching vibrations [12]. The band centered at ~ 438 cm⁻¹ can be assigned symmetric bending vibrations of Te – O – Te vertex linkages between different TeO₄ (tbp), TeO₃₊₁ polyhedra and As – O – As stretching vibrations of AsO₃ units.

4.1.7. Study by Raman spectroscopy of xMO·(100- x)[As₂O₃·PbO] vitreous systems, where MO => MnO or Fe₂O₃

The Raman spectra obtained for xMO·(100- x)[As₂O₃·PbO] vitreous systems, where MO => MnO or Fe₂O₃, with $0 \le x \le 50$ mol%, are presented in figure 4.7-8 and their structural assignments are summarized in table 4.2.

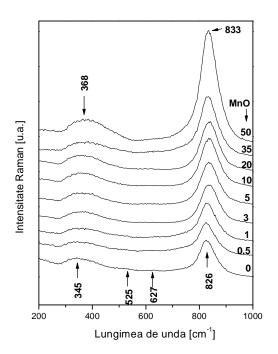


Fig. 4.7. Raman spectra of xMnO·(100- x)[As₂O₃·PbO] glasses.

The addition of MnO in the glass matrix determined an order degree of the vitreous matrix. The intensity of the bands attributed vibrations As - O - As bonds from As_2O units and bending vibrations of As - O - As from As_2O units shift to a higher wave length, that suggests an increase in the As_2O units comparative to AsO_3 units.

To the band from ~ 826 cm⁻¹ the addition of manganese ions in the glass matrix determine, for x = 0.5 mol%, an decreasing in intensity, whereupon increasing for higher concentrations. The intensity of bands centered at ~ 627 cm⁻¹ and ~ 525 cm⁻¹ for x = 0.5 % mol MnO decrease and disappear for x = 1 % mol MnO. The band of ~ 345 cm⁻¹ slowly increases with the addition of manganese ions.

To the band from ~ 826 cm⁻¹ the addition of iron ions in the glass matrix determines, for x = 0.5 mol%, an increase in intensity, a gradual decrease and disappearance for x = 35%mol Fe₂O₃ of the bands observed for the glass matrix. For higher concentrations of manganese these bands gradually decrease. The intensity of bands centered at ~ 627 cm⁻¹ and ~ 525 cm⁻¹ for x = 0.5 % mol Fe₂O₃ increase and shift to ~ 612 cm⁻¹ and ~ 501 cm⁻¹ whereas they gradually decrease and disappear for x = 10 % mol Fe₂O₃. The band of ~ 345 cm⁻¹ decreases with the addition of iron ions and disappears for x = 20 % mol Fe₂O₃.

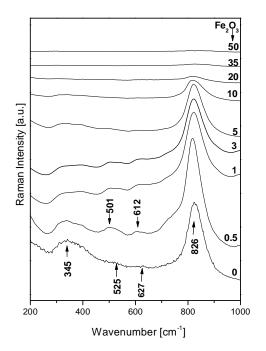


Fig. 4.8. Raman spectra of xFe₂O₃·(100- x)[As₂O₃·PbO] glasses.

With the increase of iron ions content in the As_2O_3 ·PbO glass matrix, the number of these structural units decreases due to the increasing of the disorder in the studied glasses.

4.1.8. Study by Raman spectroscopy of xMO·(100-x)[As₂O₃·TeO₂] vitreous systems, were MO => MnO or Fe₂O₃

The Raman spectra obtained for xMO·(100- x)[As₂O₃· TeO₂] vitreous systems, where MO => MnO or Fe₂O₃, with $0 \le x \le 50$ mol%, are presented in figure 4.9-10 and their structural assignments are summarized in table 4.2.

The Raman spectrum of the glass matrix suggests a structure formed from TeO_4 (tbp) and TeO_3 (tp) units where the presence of As – O – As bonds can not be excluded. The addition of manganese ions is leading to a modification of the structure proposed by the spectrum of the glass matrix. Then with the increasing of manganese content the intensity of these bands is increasing, showing that the manganese content favorises the formation of TeO_4 (tbp) and TeO_3 (tp) units.

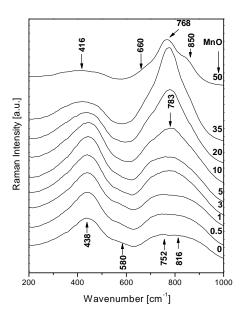


Fig. 4.9. Raman spectra of xMnO·(100- x)[As₂O₃· TeO₂] glasses.

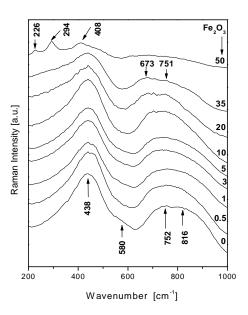


Fig. 4.10. Raman spectra of xFe₂O₃·(100- x)[As₂O₃ TeO₂] glasses.

The Raman spectra of $xFe_2O_3 \cdot (100 - x)[As_2O_3 \text{ TeO}_2]$ glasses present a significant advance to the TeO₄ (tbp), TeO₃₊₁ polyhedra and AsO₃ units represented in Raman spectrum by the ~ 438 cm⁻¹ band, comparative with the modifications determined by the addition of manganese ions in vitreous matrix. For high concentrations of iron ions in $xFe_2O_3 \cdot (100 - x)[As_2O_3 \text{ TeO}_2]$ system of glasses, the number of As-O-As linkages appear in Raman spectrum due to stretching vibrations of As-O-As and vibrations of the continuous TeO₄ (tbp) network, show a significant decrease.

4.1.9. Comparative study by Raman spectroscopy of xMnO·(100- x)[As₂O₃·BO] vitreous systems, were BO => TeO₂ or PbO

The addition of MnO to $xMnO(100-x)[As_2O_3 PbO]$ glass system do not break As - O - As bonds into characteristic units and contribute to their prominence. The addition of manganese ions to $xMnO(100-x)[As_2O_3 TeO_2]$ glass system for x=50%mol MnO, determined the splitting up of a part of As - O - As, Te-O-Te and the appearance of the TeO₄ (tbp) units against the loss of continuous TeO₄ (tbp) network was observed.

4.1.10. Comparative study by Raman spectroscopy of xFe₂O₃·(100- x)[As₂O₃·BO] vitreous systems, where BO => TeO₂ or PbO

With the increasing of iron ions content in the As₂O₃·PbO glass matrix, the number of these structural units decreased due to the increase of the disorder in the studied glasses (x > 20 % mol) and x = 50 % mol Fe₂O₃ for xFe₂O₃·(100- x)[As₂O₃·TeO₂], when iron ions brought the structural units existing in glass matrices and favoured the appearace of new bands due to bending vibrations of As – O – As.

4.2. Comparative study by FT-IR and Raman spectroscopies of xMO·(100-x)[3B₂O₃·Li₂O] glasses, where MO => MnO or Fe₂O₃ 4.2.1. Study by FT-IR spectroscopy of xMO·(100-x)[3B₂O₃·Li₂O] glasses, where MO => MnO or Fe₂O₃

The infrared absorption spectra obtained for $xMO(1-x)[3B_2O_3·Li_2O]$ glass system, where MO => MnO or Fe_2O_3 , with $0 \le x \le 50$ mol%, are presented in figure 4.11-13 and their structural assignments are summarized in table 4.3.

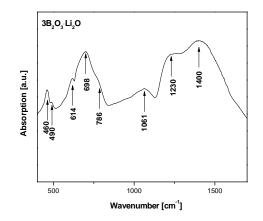


Fig. 4.11. FT – IR spectra of vitreous matrix $3B_2O_3$ ·Li₂O

Following the infrared absorption spectra obtained for $3B_2O_3 \cdot Li_2O$ vitreous matrix observed on eight bands: : ~ 1400 cm⁻¹, ~ 1230 cm⁻¹, ~ 1061 cm⁻¹, ~ 786 cm⁻¹, ~ 698 cm⁻¹, ~ 614 cm⁻¹, ~ 490 cm⁻¹ şi ~ 460 cm⁻¹. The band at ~ 1400 cm⁻¹ is assigned to asymmetric stretching vibrations from $BØ_3$ şi $BØ_2O^2$ groups (Ø represent oxygen atom bridging two boron atoms). The band at ~ 1230 cm⁻¹ is attributed to asymmetric stretching vibrations of B – O bonds from triangular ortoborate BO_3^{3-} groups [23-26].

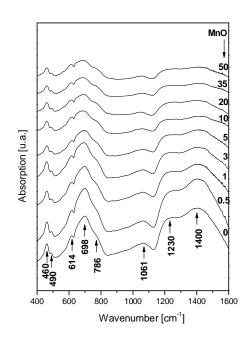


Fig. 4.12. FT – IR spectra of xMnO·(100-x)[3B₂O₃·Li₂O] glasses

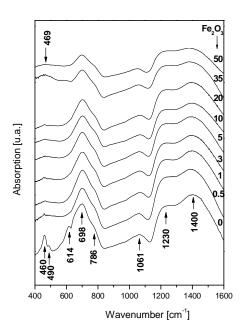


Fig. 4.13. FT – IR spectra of x Fe₂O₃·(100-x)[3B₂O₃·Li₂O] glasses 20

$Table \ 4.3. Wavenumber \ and \ the \ structural \ assignments \ of \ FT-IR \ and \ Raman \\ spectra \ of \ 3B_2O_3 \cdot Li_2O \ vitreous \ matrix$

<i>ῦ</i> [c	$2m^{-1}$]	Attribution			
FT – IR	Raman	FT – IR	Raman		
~ 1400 cm ⁻¹	~ 1400 cm ⁻¹	Asymmetric stretching vibrations from $BØ_3$ şi $BØ_2O^-$	$BØ_2O^-$ triangles linked to other borate triangular units		
		groups			
1220 -1		Asymmetric stretching			
$\sim 1230 \text{ cm}^{-1}$		vibrations of B – O bonds from			
		triangular ortoborate			
		BO ³⁻ ₃ groups			
1	1	Stretching vibrations of $B - \emptyset$	Diborate groups		
~ 1061 cm^{-1}	$\sim 1074 \text{ cm}^{-1}$	bonds in BO_4^{-} tetrahedra from			
		tri-, tetra- and penta-borate			
		groups			
	~ 902 cm ⁻¹		Vibrations of orthoborate		
	~ 902 Cm		groups		
~ 786 cm ⁻¹	~ 701 cm ⁻¹		Vibrations of rings with six		
~ 780 CIII	~ /91 Cm	– BO ₄ bonds	member to one or two BO ₄		
			tetrahedral from tri-, tetra –		
			and penta – borate units		
~ 698 cm^{-1}		Bending vibrations of B-O-B			
		bonds from pentaborate groups			
~ 614cm ⁻¹		Bending vibrations of O – B – O			
		bonds			
~ 490cm ⁻¹	~ 471 cm ⁻¹	Bending vibrations of B-O-B	Isolated diborate groups		
		bonds			
~ 460cm ⁻¹		Bending vibrations of O – B – O			
		bonds			

The band at ~ 1061 cm⁻¹ is assigned to stretching vibrations of B – Ø bonds in $BØ_4^-$ tetrahedra from tri-, tetra- and penta-borate groups and the band at ~ 786 cm⁻¹ is assigned to bending vibrations of $O_3B - O - BO_4$ bonds. The band at ~ 698 cm⁻¹ is attributed to Bending vibrations of B-O-B bonds from pentaborate groups. The bands at ~ 614 cm⁻¹ and ~ 460 cm⁻¹ are assigned to bending vibrations of O - B - O bonds, and the band at ~ 490 cm⁻¹ is assigned to bending vibrations of B-O-B bonds.

From FT-IR spectra it can be concluded that the addition of manganese or iron ions in vitreous matrices leads to the increase of the disorder degree.

Also, the structural changes that appeared in $3B_2O_3 \cdot Li_2O$ vitreous matrix with an addition MnO or Fe₂O₃ content, can be studied Ar=A₄/A₃ ratio(A₄ - relative number of BØ₄⁻ tetrahedral borate units, A₃ - relative number of triangular (BØ₃ and BØ₂O⁻) units.

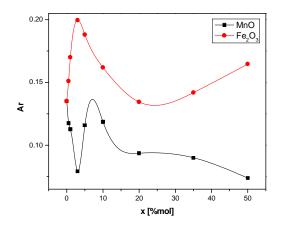


Fig. 4.14. The evolution of Ar ratio with the MnO or Fe₂O₃ content of xMnO·(100-x)[3B₂O₃·Li₂O] și xFe₂O₃·(100-x)[3B₂O₃·Li₂O]

The evolution of the value of this ratio also allows a visualization of the effects determined by the presence of the manganese or iron ions in the matrix. The value Ar << 1 indicated that BO₃ is dominant related to BO₄.

4.2.2. Study by Raman spectroscopy of xMO·(100-x)[3B₂O₃·Li₂O] glasses, where MO => MnO or Fe₂O₃

The Raman spectra obtained for $xMO(1-x)[3B_2O_3·Li_2O]$ glass system, were MO => MnO or Fe_2O_3 , with $0 \le x \le 50$ mol%, are presented in figure 4.14-16 and their structural assignments are summarized in table 4.3.

Studying the Raman spectra of $xMO \cdot (1-x)[3B_2O_3 \cdot Li_2O]$ glass system, where MO => MnO or Fe_2O_3 , it has been observed that the absorption bands presented in vitreous matrix gradually decreased with an increase in the manganese or iron content, in both glasses system, which suggested an increase of the disorder in these vitreous systems.

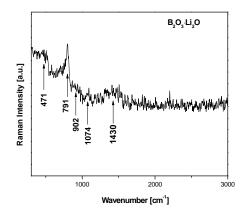


Fig. 4.15. Raman spectra of vitreous matrix 3B₂O₃·Li₂O

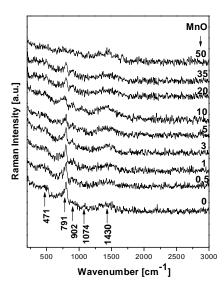


Fig. 4.16. Raman spectra of xMnO·(100-x)[3B₂O₃·Li₂O] glasses

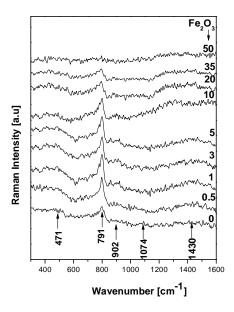


Fig. 4.17. Raman spectra of x Fe₂O₃·(100-x)[3B₂O₃·Li₂O] glasses

4.3. Comparative study by electron paramagnetic resonance (EPR) and magnetic susceptibility measurements of xMO·(100-x)[As₂O₃·TeO₂], xMO·(100-x)

 $[As_2O_3 \cdot PbO]$ and $xMO \cdot (100 \cdot x)[3B_2O_3 \cdot Li_2O]$ glasses, where MO => MnO or Fe_2O_3

4.3.1. Comparative study by electron paramagnetic resonance (EPR) of xMnO·(100-x)[As₂O₃·BO], where BO => TeO₂ or PbO and xMnO·(100-x) [3B₂O₃·Li₂O] glasses

In order to obtain more information regarding the behavior of manganese ions in different vitreous matrices [30,31], xMnO·(100-x)[As₂O₃·BO], where BO => TeO₂ or PbO and xMnO·(100-x) [3B₂O₃·Li₂O] glasses were investigated by EPR spectroscopy on a large compositional range, with $0 \le x \le 50$ mol%. The EPR spectra of Mn²⁺ for xMnO·(100-x)[As₂O₃·PbO] (S1), xMnO·(100-x)[As₂O₃·TeO₂] (S3) and xMnO·(100-x)[3B₂O₃·Li₂O] (S5) are presented in figure 4.18-20. These spectra present two absorption lines centered at g_{ef} \approx 4,3 and g_{ef} \approx 2,0 in all systems.

The absorption line centered at $g_{ef} \approx 4,3$ and is due to isolated Mn^{2+} ions [32]. The absorption line centered at $g_{ef} \approx 2,0$ may be attributed to Mn^{2+} species interacting by magnetic coupling[33], dipolar and/ or super exchange, the last ones forming magnetic clusters [32].

The evolution of the spectra is easier to follow considering the dependence of concentration on the EPR parameters, the line - intensity (obtained as an integral of the area under the corresponding EPR signal), J and the line - width, ΔB . The evolution of J and AB reflects the structural transformations which appear in the glass matrices due to the increase of manganese ions content.

The intensity and the line - width of the resonance line from $g_{ef} \approx 4,3$ for all investigated systems is represented in figure 4.21(a,b).

The intensity of the resonance line from $g_{ef} \approx 4,3$ for all investigated glass systems decreases due to the decrease of the isolated manganese ions numbers. The line - width from $g_{ef} \approx 4,3$ for (S1) and (S3) systems decreases due to decrease of Mn²⁺ number and to the structural disorder in glasses with the increase of MnO content. For (S5) glass system the line - width of the resonance line there is an increase in all compositional domains.

The intensity of the resonance line from $g_{ef} \approx 2,0$ can be observed as increasing up to x=14% mol for (S1) system, up to x=8%mol for (S3) system and 20 %mol for (S5) system. The line - width of the resonance line from $g_{ef} \approx 2,0$ for (S1) and (S3) systems increases up to 3 % mol and 5%mol for (S5) due to Mn^{2+} species interacting by magnetic coupling dipole- dipole as the main broadening mechanism. Over this concentration line – the width of the resonance line from $g_{ef} \approx 2,0$ is straight, and it can be assumed that the manganese ions begin to manifest super exchange magnetic interactions.

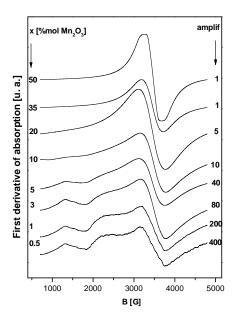


Fig. 4.18. The EPR spectra of Mn²⁺ for xMnO·(100-x)[As₂O₃·PbO] (S1) glass system

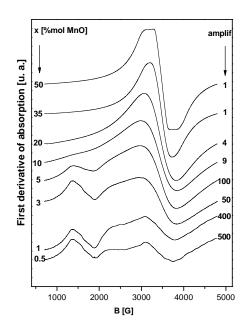


Fig. 4.19. The EPR spectra of Mn²⁺ for xMnO·(100-x)[As₂O₃·TeO₂] (S3) glass system.

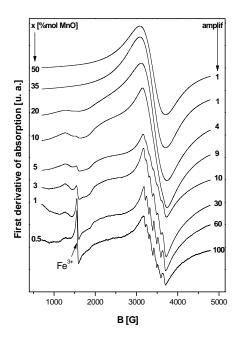


Fig. 4.20. The EPR spectra of Mn²⁺ for xMnO·(100-x)[3B₂O₃·Li₂O] (S5) glass system

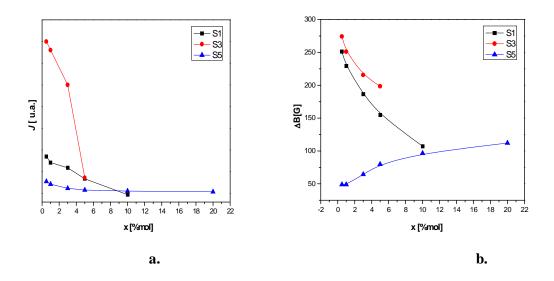


Fig. 4.21. The dependence on MnO content of the intensity (a) and width(b) of resonance line at g_{ef}≈ 4,3 for xMnO·(100-x)[As₂O₃·PbO] (S1), xMnO·(100-x)[As₂O₃·TeO₂] (S3) şi xMnO·(100-x)[3B₂O₃·Li₂O] (S5) system glasses

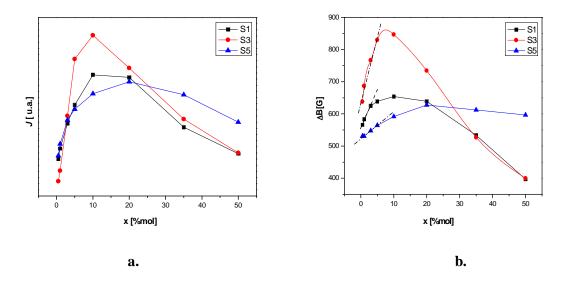


Fig. 4.22. The dependence on MnO content of the intensity (a) and width b) of resonance line at $g_{ef} \approx 2,0$ for xMnO·(100-x)[As₂O₃·PbO] (S1), xMnO·(100-x)[As₂O₃·TeO₂] (S3) and xMnO·(100-x)[3B₂O₃·Li₂O] (S5) system glasses

Following figure 4.18-20 it can be observed that (S1) and (S3) have had a similar evolution, but (S5) system presented for $0.5 \le x \le 5$ %mol (fig.4.20) has a hyperfine sextet[33] at $g_{ef} \approx 2.0$, due to Mn^{2+} isolated ions.

4.3.2. Comparative study by magnetic susceptibility measurements of xMnO·(100-x)[As₂O₃·BO], where BO => TeO₂ or PbO and xMnO·(100-x)[3B₂O₃·Li₂O] glasses

The magnetic susceptibility data are in good agreement with the EPR result. The magnetic behavior of the glasses from $xMnO(100-x)[As_2O_3 \cdot BO]$, where $BO => TeO_2$ (S3) or PbO (S1) and $xMnO(100-x)[3B_2O_3 \cdot Li_2O]$ (S5) glasses for $3 \le x \le 50$ mol% MnO are presented in figure 4.23 - 25.

For (S1) and (S3) system glasses up to $x \le 3$ %mol and $x \le 5$ %mol (S5) glass system, the temperature dependence of the reciprocal magnetic susceptibility obeys a Curie law. In this concentration range the manganese ions are predominantly isolated or/and participate in dipole-dipole interractions. At higher concentrations the reciprocal magnetic susceptibility obeys a Curie-Weiss law with negative paramagnetic Curie temperature: (θ_p) characteristic to antiferromagnetic coupled ions by means of super exchange interractions(fig. 4.26). This behavior of manganese ions is confirmed by the shape of the EPR spectra.

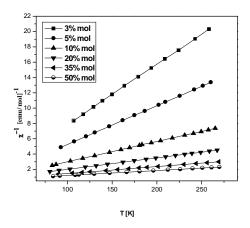


Fig.4.23. The temperature dependence of χ^{-1} for xMnO·(100-x)[As₂O₃·PbO],

(S1), $3 \le x \le 50$ % mol

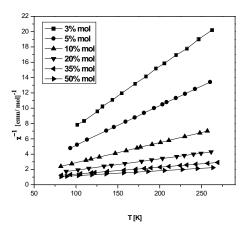


Fig.4.24. The temperature dependence of χ^{-1} for xMnO·(100-x)[As₂O₃·TeO₂], (S1), $3 \le x \le 50$ % mol

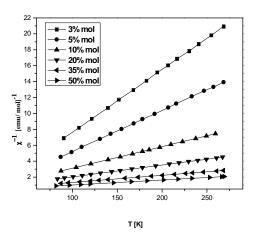


Fig.4.25. The temperature dependence of χ^{-1} for xMnO·(100-x)[3B_2O_3·Li_2O] (S5), $3 \le x \le 50$ %mol

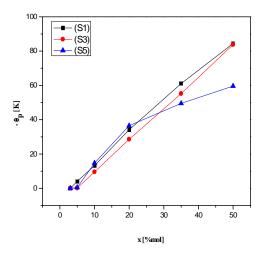


Fig.4.26.Concentration dependence of (θ_p) for sistemele xMnO·(100-x)[As₂O₃·PbO] (S1), xMnO·(100-x)[As₂O₃·TeO₂] (S3) şi xMnO·(100-x)[3B₂O₃·Li₂O] (S5).

Using the representation $1/\chi = f(T)$ the molar Curie constants, C_M , the effective magnetic moments and molar fraction of $Mn^{2+}(x_1)$ şi $Mn^{3+}(x_2)$ ions were calculated and represented in tables 4.4 -6.

Tabel 4.4. Molar Curie constants, effective magnetic moments and molar fraction of
$Mn^{2+}(x_1)$ și $Mn^{3+}(x_2)$ ions in xMnO·(100-x)[As ₂ O ₃ ·PbO] (S1) glasses

X	$C_M \times 10^2$	μ_{ef}	X ₁	X2
[%mol MnO ₂]	[emu/ mol]	[µ _B]	$[\% mol Mn^{2+}O_2]$	$[\% mol Mn^{3+}O_2]$
3	12,72	5,82	2,68	0,32
5	19,79	5,62	3,43	1,57
10	37,77	5,49	5,55	4,45
20	66,88	5,17	4,92	15,08
35	122,39	5,02	3,77	31,23
50	154,27	4,96	2,68	47,32

 $\begin{array}{l} \mbox{Tabel 4.5. Molar Curie constants, effective magnetic moments and molar fraction of $Mn^{2+}(x_1)$ is $Mn^{3+}(x_2)$ ions in $xMnO(100-x)[As_2O_3(TeO_2)]$ (S3)glasses } \end{array} \end{array}$

X	$C_M \times 10^2$	μ_{ef}	x ₁	X2
[%mol MnO ₂]	[emu/ mol]	[μ _B]	$[\% mol Mn^{2+}O_2]$	$[\% mol Mn^{3+}O_2]$
3	12,97	5,88	2,87	0,13
5	19,38	5,56	3,12	1,88
10	38,04	5,51	5,75	4,25
20	67,66	5,2	5,49	14,51
35	111,48	5,04	4,41	30,59
50	155,28	4,98	3,58	46,42

X	$C_M \times 10^2$	μ_{ef}	x ₁	X2
[%mol MnO ₂]	[emu/ mol]	[µ _B]	$[\% mol Mn^{2+}O_2]$	[%mol Mn ³⁺ O ₂]
3	12,7	5,81	2,65	0,35
5	19,19	5,54	3,03	1,97
10	36,81	5,42	4,86	5,14
20	72,15	5,19	5,3	14,7
35	112,1	5,06	5,05	29,95
50	158,98	5,01	4,93	45,07

Tabel 4.6. Molar Curie constants, effective magnetic moments and molar fraction of Mn²⁺(x₁) și Mn³⁺(x₂) ions in xMnO·(100-x)[3B₂O₃·Li₂O] (S5) glasses.

4.3.3. Comparative study by electron paramagnetic resonance (EPR) of xFe₂O₃·(100-x)[As₂O₃·BO], where BO => TeO₂ or PbO and x Fe₂O₃·(100-x) [3B₂O₃·Li₂O] glasses

In order to obtain more information regarding the behavior of iron ions in $xFe_2O_3 \cdot (100-x)[As_2O_3 \cdot BO]$, where BO => TeO₂ (S4) or PbO (S2) şi $xFe_2O_3 \cdot (100-x)[3B_2O_3 \cdot Li_2O]$, glasses were investigated by EPR spectroscopy on a large compositional range, with 0,5 $\leq x \leq 35$ mol% Fe₂O₃. The EPR spectra of Fe³⁺ in $xFe_2O_3 \cdot (100-x)[As_2O_3 \cdot PbO]$ (S2), $xFe_2O_3 \cdot (100-x)[As_2O_3 \cdot TeO_2]$ (S4) and $xFe_2O_3 \cdot (1-x)[3B_2O_3 \cdot Li_2O]$ (S6) are presented in figure 4.27-29. These spectra present two absorption lines centered at $g_{ef} \approx 4,3$ and $g_{ef} \approx 2,0$ up to $x \leq 35\%$ mol Fe₂O₃ systems studied.

The absorption line centered at $g_{ef} \approx 4,3$ is due to isolated Fe³⁺ ions [34-36]. The absorption line centered at $g_{ef} \approx 2,0$ may be attributed to Fe³⁺ species interacting by magnetic coupling, dipolar and/ or super exchange, the last ones forming magnetic clusters [37-41].

The evolution of the spectra is easier to follow considering the dependence of concentration on the EPR parameters, the line - intensity (obtained as an integral of the area under the corresponding EPR signal), J and the line - width, ΔB . The evolution of J and AB reflects the structural transformations which appear in the glass matrices due to the increasing of iron ions content.

The intensity and the line - width of the resonance line from $g_{ef} \approx 4,3$ for all investigated systems is represented in figure 4.30 (a, b). The intensity of the resonance line from $g_{ef} \approx 4,3$ increases for (S2) and (S4) up to 3 % mol and up to 5 % mol for (S6) then decreases for all systems investigated. A line broadening up to 3 % mol from $g_{ef} \approx 4,3$ for these systems can be observed due to increasing of Fe³⁺ isolated number of ions.

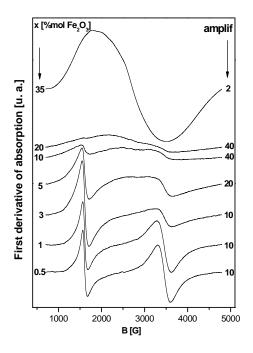


Fig.4.27. The EPR spectra of Fe³⁺ for xFe₂O₃·(100-x)[As₂O₃·PbO] (S2) glass system

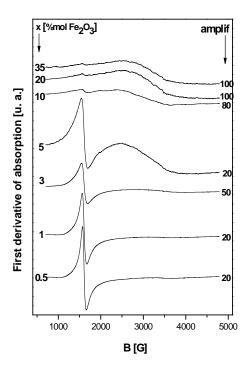


Fig.4.28. The EPR spectra of Fe³⁺ for xFe₂O₃·(100-x)[As₂O₃·TeO₂] (S4) glass system

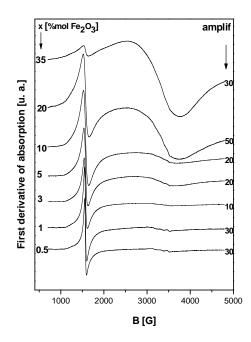


Fig.4.29. The EPR spectra of Fe³⁺ for xFe₂O₃·(100-x)[3B₂O₃·Li₂O] (S6) glass system

The intensity and the line - width of the resonance line from $g_{ef} \approx 4,3$ for all investigated systems is represented in figure 4.30 (a, b). The resonance line from $g_{ef} \approx 2,0$ is attributed to Fe³⁺ ions that interact by magnetic coupling, dipolar or/and super exchange interractions can be observed in all compositional domains.

The intensity of the resonance line from $g_{ef} \approx 2,0$ can be observed as increasing up to x=35% mol for (S2) and (S6) system and 10 %mol for (S4) system with increased iron ions content. The line - width of the resonance line from $g_{ef} \approx 2,0$ increased up to 10%mol for (S2) system, and up to 3 % mol for (S4) and 5%mol for (S6) due to Fe³⁺species interacting by magnetic coupling dipole-dipole. Above this concentration line - width of the resonance line from $g_{ef} \approx 2,0$ can be attributed to dipolar and super exchange magnetic interactions. The $\Delta B = f(x)$ dependence reflects the competion between the broadening mechanism (dipole – dipole interactions, the interactions between ions in different valence state and structural disorder) and the narrowing ones (super exchange interactions).

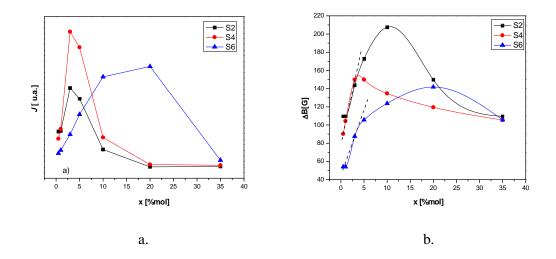


Fig. 4.30. The dependence on Fe₂O₃ content of the intensity (a) and width (b) of resonance line at $g_{ef} \approx 4,3$ for $xFe_2O_3 \cdot (100-x)[As_2O_3 \cdot PbO]$ (S2), $xFe_2O_3 \cdot (100-x)[As_2O_3 \cdot TeO_2]$ (S4) and $xFe_2O_3 \cdot (100-x)[3B_2O_3 \cdot Li_2O]$ (S6) glasses

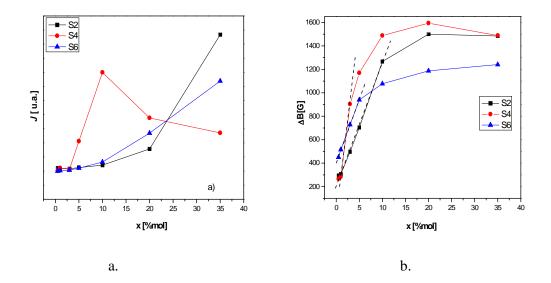


Fig. 4.31. The dependence on Fe₂O₃ content of the intensity (a) and width (b) of resonance line at g_{ef} ≈ 2,0 for xFe₂O₃·(100-x)[As₂O₃·PbO] (S2), xFe₂O₃·(100-x)[As₂O₃·TeO₂] (S4) and xFe₂O₃·(100-x)[3B₂O₃·Li₂O] (S6) glasses

4.3.4. Comparative study by magnetic susceptibility measurements of xFe₂O₃·(100-x)[As₂O₃·BO], where BO => TeO₂ or PbO and xFe₂O₃·(100-x) [3B₂O₃·Li₂O] glasses

The magnetic susceptibility data are in good agreement with the EPR result. The magnetic behavior of the glasses $xFe_2O_3 \cdot (100-x)[As_2O_3 \cdot PbO]$ (S2), $xFe_2O_3 \cdot (100-x)[As_2O_3 \cdot TeO_2]$ (S4) and $xFe_2O_3 \cdot (100-x)[3B_2O_3 \cdot Li_2O]$ (S6) for $3 \le x \le 50$ % mol Fe₂O₃ are presented in figure 4.32 - 34.

For (S2) glass systems up to $x \le 10$ %mol, $x \le 3$ %mol (S5) glass systems and $x \le 5$ %mol (S6) glass systems the temperature dependence of the reciprocal magnetic susceptibility obeys a Curie law. In this concentration range the iron ions are predominantly isolated or/and participate in dipole-dipole interactions. At higher concentrations the reciprocal magnetic susceptibility obeys a Curie-Weiss law with negative paramagnetic Curie temperature: (θ_p) characteristic to antiferromagnetic coupled ions by means of super exchange interactions(fig. 4.35). This behavior of iron ions is confirmed by the shape of the EPR spectra.

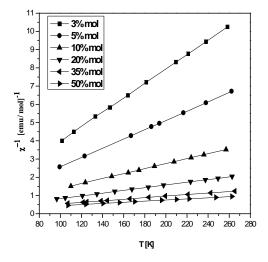


Fig.4.32. The temperature dependence of χ^{-1} for xFe₂O₃·(100-x)[As₂O₃·PbO] (S2), 3 ≤ x ≤ 50 % mol

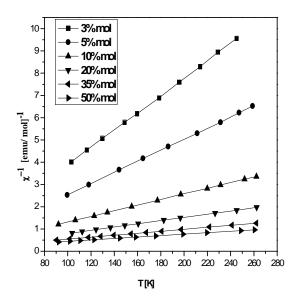


Fig.4.33. The temperature dependence of χ^{-1} for $xFe_2O_3 \cdot (100-x)[As_2O_3 \cdot TeO_2]$ (S4),

 $3 \le x \le 50$ % mol

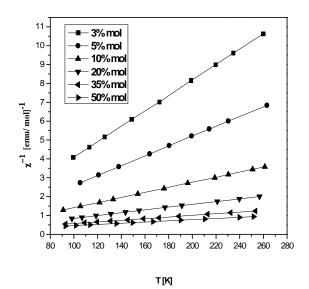
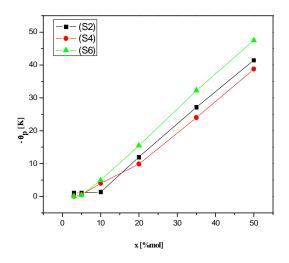


Fig.4.34. The temperature dependence of χ^{-1} for xFe₂O₃·(100-x)[3B₂O₃·Li₂O] (S6), 3 ≤ x ≤ 50 % mol



 $\label{eq:states} \begin{array}{l} Fig.4.35. Concentration dependence of (θ_p) for sistemele xFe_2O_3(100-x)[As_2O_3$PbO]$ (S2), xFe_2O_3(100-x)[As_2O_3$TeO_2]$ (S4) and xFe_2O_3(100-x)[3B_2O_3$Li_2O]$ (S6), $for $3 \leq x \leq 50 \mbox{ } \mbox{mol } Fe_2O_3$ (S6), $for $3 \leq x \leq 50 \mbox{mol } Fe_2O_3$ (S6), $for $3 \leq x < 50 \mbox{mol } Fe_2O_3$ (S6), $for $3 \leq x < 50 \mbox{mol } Fe_2O_3$ (S6), $for $3 \leq x < 50 \mbox{mol } Fe_2O_3$ (S6), $for $3 \leq x < 50 \mbox{mol } Fe_2O_3$ (S6), $for $3 \leq x < 50 \mbox{mol } Fe_2O_3$ (S6), $for $3 \leq x < 50 \mbox{mol } Fe_2O_3$ ($

Using the representation $1/\chi = f(T)$ the molar Curie constants, C_M , the effective magnetic moments and molar fraction of $Fe^{3+}(x_1)$ și $Fe^{2+}(x_2)$ ions were calculated and represented in tables 4.7-9.

X	$C_M \times 10^2$	μ_{ef}	X ₁	X2
[%mol Fe ₂ O ₃]	[emu/ mol]	[µ _B]	$[\% mol Fe_2^{3+}O_3]$	$[\% mol Fe_2^{2+}O_3]$
3	25,26	5,8	2,61	0,39
5	39,3	5,6	3,32	1,68
10	73,26	5,41	4,76	5,24
20	134,59	5,19	5,3	14,7
35	233,64	5,16	8,29	26,71
50	321,54	5,07	7,67	42,33

Table 4.7. Molar Curie constants, effective magnetic moments and molar fraction of $Fe^{3+}(x_1)$ și $Fe^{2+}(x_2)$ ions in $xFe_2O_3 \cdot (100-x)[As_2O_3 \cdot PbO]$ glasses (S2)

Table 4.8. Molar Curie constants, effective magnetic moments and molar fraction of $Fe^{3+}(x_1)$ și $Fe^{2+}(x_2)$ ions in $xFe_2O_3 \cdot (100-x)[As_2O_3 \cdot TeO_2]$ glasses (S4)

Х	$C_M \times 10^2$	μ_{ef}	X1	X2
[%mol Fe ₂ O ₃]	[emu/ mol]	$[\mu_B]$	$[\% mol Fe_2^{3+}O_3]$	$[\% mol Fe_2^{2+}O_3]$
3	25,71	5,85	2,77	0,23
5	40,16	5,66	3,63	1,37
10	79,11	5,62	6,86	3,14
20	149,92	5,21	6,43	9,29
35	226,76	5,09	6,01	28,99
50	309,59	4,97	3,13	46,87

Х	$C_M \times 10^2$	μ_{ef}	X1	X2
[%mol Fe ₂ O ₃]	[emu/ mol]	$[\mu_B]$	$[\% mol Fe_2^{3+}O_3]$	$[\% mol Fe_2^{2+}O_3]$
3	25,42	5,82	2,68	0,32
5	39,43	5,61	3,38	1,62
10	74,23	5,44	5,06	4,94
20	135,68	5,20	5,49	14,51
35	230,94	5,13	7,31	27,69
50	317,46	5,03	5,84	44,16

Table 4.9. Molar Curie constants, effective magnetic moments and molar fraction of Fe³⁺(x₁) și Fe²⁺(x₂) ions in xFe₂O₃·(100-x)[3B₂O₃·Li₂O] glasses (S6)

The results allowed to consider that the magnetic properties of the studied glasses can be explained through magnetic interactions between $Fe^{3+} - Fe^{3+}$, $Fe^{3+} - Fe^{2+}$ and $Fe^{2+} - Fe^{2+}$ that explain the dependences of RPE parameters concentration.

SELECTED CONCLUSIONS

In this thesis, $xMO \cdot (100-x)[As_2O_3 \cdot TeO_2]$, $xMO \cdot (100-x) [As_2O_3 \cdot PbO]$ and $xMO \cdot (100-x)[3B_2O_3 \cdot Li_2O]$ glasses, where MO => MnO or Fe_2O_3 , were comparatively studied in $0 \le x \le 50$ %mol. The structural modifications of investigated glass matrices were comparatively followed in function of the manganese or iron ions on the investigated glass matrices were followed.

The study by FT-IR spectroscopy shows that the evolution of the spectra of xMnO(100-x) [As₂O₃·PbO] with the addition and the increasing of the MnO content suggest that the manganese ions break up a part of As – O – As, Pb – O – Pb and probably of As – O – Pb bonds. The addition of MnO in the As₂O₃·PbO glass matrix is leading to a disordering of the glasses structure with the increasing of the MnO content.

The evolution of the absorption bands for FT - IR spectrum of $xFe_2O_3 \cdot (1-x)$ [As₂O₃·PbO] glasses are determined by the addition of Fe₂O₃ content. It can be observed to shift the bands assigned As – O – As vibrations from AsO₃ units due to changing to length As – O – As bonds and interne angles of AsO₃ units, disappears PbO₄ units and appears FeO₆ units The addition of Fe₂O₃ in the As₂O₃·PbO glass matrix determined a disordering of the glasses structure with the increasing of the Fe₂O₃ content.

From FT-IR absorption spectra of the investigated $xMO \cdot (100-x)[As_2O_3 \cdot TeO_2]$ glasses system where MO => MnO or Fe_2O_3 , it can be observed that the addition and increasing of manganese ions concluded: the contraction of As-O-As bonds from totally symmetric stretching vibrations of AsO₃ units; for $x \ge 10$ % mol there can be an increase in number of As-O-As bonds from AsO₃ units or/and a length these linkage to doubly degenerate stretching vibrations of AsO₃ units and a stretching vibration mode of TeO₃ tp with NBO. It can be remarked from the FT-IR absorption spectra, the gradual transformation of symmetric stretching vibration mode of Te-O bonds from TeO₃ units in TeO₄ (tbp) units and the presence of MnO₂ units. Once, with the addition of iron ions, the xFe₂O₃·(100- x)[As₂O₃·TeO₂] glasses behave like xMnO·(100-x)[As₂O₃·TeO₂] glasses, but in this system the appearance of FeO₆ units can be seen.

From FT-IR spectra in $xMO(100-x)[3B_2O_3\cdot Li_2O]$ glasses, where MO => MnO or Fe_2O_3 can be concluded that the addition of manganese or iron ions in vitreous matrices leads to the increase of the disorder degree.

The study by Raman spectroscopy shows that the evolution of the spectra the addition of MnO in the $[As_2O_3 \cdot PbO]$ glass matrix determined an order degree of the vitreous matrix. The intensity of the bands attributed vibrations As – O – As bonds from As₂O units and bending vibrations of As – O – As from As₂O units shift to a higher wave length, that suggests an increase in the As₂O units comparative to AsO₃ units. With the increase of iron ions content in the As₂O₃·PbO glass matrix, the number of these structural units decreases due to the increasing of the disorder in the studied glasses.

With the increasing of manganese content in As_2O_3 . TeO₂ glass matrix, the intensity of these bands is increasing, showing that the manganese content favorites the formation of TeO₄ (tbp) and TeO₃ (tp) units. For high concentrations of iron ions in xFe₂O₃·(100x)[As₂O₃ TeO₂] system of glasses, the number of As-O-As linkages appear in Raman spectrum due to stretching vibrations of As-O-As and vibrations of the continuous TeO₄ (tbp) network, show a significant decrease.

Studying the Raman spectra of $xMO(1-x)[3B_2O_3\cdot Li_2O]$ glass system, where MO => MnO or Fe_2O_3 , it has been observed that the absorption bands presented in vitreous matrix gradually decreased with an increase in the manganese or iron content, in both glasses system, which suggested an increase of the disorder in these vitreous systems.

The EPR spectra presents the behavior of manganese or iron ions in different vitreous matrices, $xMO\cdot(100-x)[As_2O_3\cdot BO]$, where MO=>MnO or Fe_2O_3 and $BO => TeO_2$ or PbO and $xMO\cdot(100-x)$ [$3B_2O_3\cdot Li_2O$]. Glasses were investigated by EPR spectroscopy on a large compositional range, with $0 \le x \le 50$ mol%. The absorption line centered at $g_{ef} \approx 4,3$ and is due to isolated Mn^{2+} ions. The absorption line centered at $g_{ef} \approx 2,0$ may be attributed to Mn^{2+} species interacting by magnetic coupling, dipolar and/ or super exchange, the last ones forming magnetic clusters. The EPR spectra of Fe^{3+} in $xFe_2O_3\cdot(100-x)[As_2O_3\cdot PbO]$, $xFe_2O_3\cdot(100-x)[As_2O_3\cdot TeO_2]$ and $xFe_2O_3\cdot(1-x)[3B_2O_3\cdot Li_2O]$ are presented two absorption lines centered at $g_{ef} \approx 4,3$ and $g_{ef} \approx 2,0$ up to $x \le 35\%$ mol Fe_2O_3 systems studied.

The absorption line centered at $g_{ef} \approx 4,3$ is due to isolated Fe^{3+} ions . The absorption line centered at $g_{ef} \approx 2,0$ may be attributed to Fe^{3+} species interacting by magnetic coupling, dipolar and/ or super exchange, the last ones forming magnetic clusters.

The evolution of the spectra is easier to follow considering the dependence of concentration on the EPR parameters, the line - intensity (obtained as an integral of the area under the corresponding EPR signal), J and the line - width, ΔB . The evolution of J and AB reflects the structural transformations which appear in the glass matrices due to the increase of manganese ions content.

The magnetic susceptibility data are in good agreement with the EPR result. The magnetic behavior of the glasses from $xMnO(100-x)[As_2O_3 \cdot BO]$, where $BO => TeO_2$ (S3) or PbO (S1) and $xMnO(100-x)[3B_2O_3 \cdot Li_2O]$ (S5) glasses for $3 \le x \le 50$ mol% MnO.

For (S1) and (S3) system glasses up to $x \le 3$ % mol and $x \le 5$ % mol (S5) glass system, the temperature dependence of the reciprocal magnetic susceptibility obeys a Curie law. In this concentration range the manganese ions are predominantly isolated or/and participate in dipole-dipole interractions. At higher concentrations the reciprocal magnetic susceptibility obeys a Curie-Weiss law with negative paramagnetic Curie temperature: (θ_p) characteristic to antiferromagnetic coupled ions by means of super exchange interractions. This behavior of manganese ions is confirmed by the shape of the EPR spectra. Using the representation $1/\chi = f(T)$ the molar Curie constants, C_M , the effective magnetic moments and molar fraction of $Mn^{2+}(x_1)$ şi $Mn^{3+}(x_2)$ ions were calculated. The magnetic susceptibility data are in good agreement with the EPR result.

The magnetic behavior of the glasses $xFe_2O_3 \cdot (100-x)[As_2O_3 \cdot PbO]$ (S2), $xFe_2O_3 \cdot (100-x)[As_2O_3 \cdot TeO_2]$ (S4) and $xFe_2O_3 \cdot (100-x)[3B_2O_3 \cdot Li_2O]$ (S6) for $3 \le x \le 50$ %mol Fe_2O_3 are presented. For (S2) glass systems up to $x \le 10$ %mol, $x \le 3$ %mol (S5) glass systems and $x \le 5$ %mol (S6) glass systems the temperature dependence of the reciprocal magnetic susceptibility obeys a Curie law. In this concentration range the iron ions are predominantly isolated or/and participate in dipole-dipole interactions. At higher concentrations the reciprocal magnetic susceptibility obeys a Curie-Weiss law with negative paramagnetic Curie temperature: (θ_p) characteristic to antiferromagnetic coupled ions by means of super exchange interactions(fig. 4.35). This behavior of iron ions is confirmed by the shape of the EPR spectra. Using the representation $1/\chi = f(T)$ the molar Curie constants, C_M , the effective magnetic moments and molar fraction of $Fe^{3+}(x_1)$ şi $Fe^{2+}(x_2)$ ions were calculated.

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